

Magnetar Central Engine Powering the Energetic GRB 210610B?

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Abstract

The bright GRB 210610B was discovered simultaneously by *Fermi* and *Swift* missions at redshift 1.13. We utilized broadband *Fermi*-GBM observations to perform a detailed prompt emission spectral analysis and to understand the radiation physics of the burst. Our analysis displayed that the low energy spectral index (α_{pt}) exceeds boundaries expected from the typical synchrotron emission spectrum (-1.5 , -0.67), suggesting additional emission signature. We added an additional thermal model with the typical Band or CPL function and found that CPL + BB function is better fitting to the data, suggesting a hybrid jet composition for the burst. Further, we found that the beaming corrected energy ($E_{\gamma,\theta_j} = 1.06 \times 10^{51}$ erg) of the burst is less than the total energy budget of the magnetar. Additionally, the X-ray afterglow light curve of this burst exhibits achromatic plateaus, adding another layer of complexity to the explosion’s behavior. Interestingly, we noted that the X-ray energy release during the plateau phase ($E_{\text{X,iso}} = 1.94 \times 10^{51}$ erg) is also less than the total energy budget of the magnetar. Our results indicate the possibility that a magnetar could be the central engine for this burst.

Keywords: Gamma-ray burst, Synchrotron, Thermal, Magnetar

1. Introduction

Gamma ray bursts (GRBs) are the most energetic explosions in the universe. GRBs emit electromagnetic radiation in two successive phases. The initial “prompt emission phase” is a complex function of time, energy, and polarization, spanning a broad range of frequencies from radio to gamma rays up to TeV energies (MAGIC Collaboration et al., 2019), that poses significant challenges for relating it to known physical emission processes. The synchrotron emission mechanism can explain the observed spectral features in some GRBs. However, some bursts exhibit a spectral component that appears to be inconsistent with synchrotron emission, such as the low energy spectral index α_{pt} of the Band function not always remaining within the limits

(-1.5 , -0.67) known as the synchrotron line of death (Preece et al., 1998). GRBs deviating from these limits can be explained by adding a thermal component to the spectrum, indicating the hybrid jet composition for these bursts (Kumar and Zhang, 2015; Pe’er, 2015). Following this, the “afterglow emission phase” of a GRB is generated through the synchrotron radiation from the electrons that are accelerated in the external shock formed by the interaction of the GRB ejecta with the surrounding medium. Theoretical models based on this mechanism have been able to reproduce the observed afterglow light curves and spectral properties with remarkable accuracy. However, some observed deviations from the expected afterglow emission, such as achromatic plateaus or bumps, continue to challenge the synchrotron model and may indicate the presence of more complex emission processes. In recent years, a handful of GRBs have been identified that exhibit a plateau phase in their afterglow light curves. The plateau phase is believed to be a result of energy injection into the external shock, which maintains the shock’s constant energy over a longer timescale. The injection can be achieved through different mechanisms; the most plausible scenario is the magnetar central engine, where the energy injection from a magnetar can cause the external plateau in the afterglow light curves (Stratta et al., 2018). The cosmological constants chosen for this article are Hubble parameter $H_0 = 71 \text{ km s}^{-1} \text{ Mpc}^{-1}$, density parameters $\Omega_\Lambda = 0.73$, and $\Omega_m = 0.27$.

2. Data Analysis and Results

For the present analysis, we have used the publicly available data of GRB 210610B, which was detected by *Fermi* on 2021-06-10 at 19:51:05.05 UT (Malacaria et al., 2021) as well as by *Swift* and several other space- and ground-based telescopes during the prompt and afterglow phase.

2.1. Prompt emission

The Gamma-Ray Burst Monitor (GBM, Meegan et al. 2009) onboard *Fermi* is specifically designed to detect the prompt emissions of GRBs with excellent temporal and spectral resolution. To analyse the GBM data, we utilized the Multi-Mission Maximum Likelihood (3ML, Vianello et al. 2015) framework. We downloaded the GBM data from the *Fermi* Science Support Center (FSSC) and utilized time-tagged event files from three NaI detectors and one BGO detector with the minimum deviation from the direction of the burst. The time-integrated and time-resolved spectra were then extracted using the *gtburst* package. To create the time-resolved spectrum, time slicing was performed by utilizing the Bayesian block binning method with a false alarm probability of 0.01. The extracted spectrum was loaded into 3ML utilizing the GBM plugin, and we employed various inbuilt empirical models, such as Band and Cutoff powerlaw (CPL), along with physical models like blackbody (BB) and physical synchrotron (Burgess et al., 2020), and their combinations for spectral fitting. We adopted the Bayesian method to fit the model to the data and evaluated the goodness-of-fit using the Deviance Information Criterion (DIC) statistical test (Spiegelhalter et al., 2002). The model with the lowest DIC value was considered the best fit. For further details about the data analysis and model comparison, please refer to Ror et al. (2023).

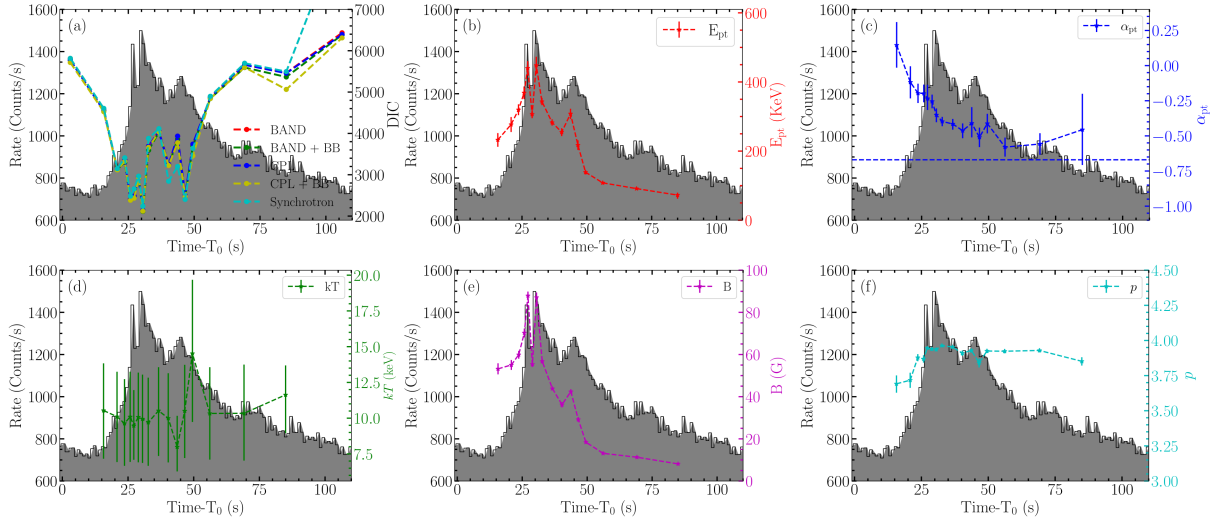


Figure 1: Results of the time-resolved spectral analysis of GRB 210610B. Panel (a) represents the DIC comparison of various models used to fit the time-resolved prompt emission spectrum. Panels (b), (c), and (d) represent the evolution of peak energy E_{pt} , low energy spectral index α_{pt} , and temperature (kT) of the fireball from the best-fit model (CPL+BB) along with the prompt emission light curve in the background. Panels (e) and (f) represent the evolution of magnetic field strength B and electron energy distribution index p from the physical synchrotron model.

Time-integrated spectrum analysis shows that the Band+BB function provided the most accurate fit with the lowest DIC value. The obtained peak energy $E_{\text{pt}} = 284^{+5.50}_{-5.56}$ keV, isotropic energy $E_{\gamma,\text{iso}} = 4.31 \times 10^{53}$ erg, and peak luminosity $L_{\gamma,\text{iso}} = 7.08 \times 10^{52}$ erg s $^{-1}$, for GRB 210610B are consistent with the well-studied Amati (Amati, 2006) and Yonetoku (Yonetoku et al., 2004) correlations. The results of the time-resolved spectral analysis are shown in Fig. 1. A comparison of DIC values indicated that most of the bins are best fitted by the CPL+BB model, while the physical synchrotron model is favoured by some of the bins. Spectral parameters obtained from the best-fit models in the time-resolved analysis show the following evolution pattern: The peak energy (E_{pt}) and the magnetic field strength (B) are found to track the observed flux. Furthermore, the low energy spectral index (α_{pt}) is found crossing the synchrotron line of death and showing the hard to soft evolution, while the electron energy distribution index (p) has remained almost constant throughout the burst.

2.2. X-ray afterglow

To study the X-ray afterglow, we downloaded the *Swift*-XRT observations (Page et al., 2021; Gropp et al., 2021) from the archive at the www.swift.ac.uk. To fit the *Swift*-XRT light curve, we employed broken power-law models with one, two, and three breaks. The fitting was performed using the QDP package, and the χ^2 statistic was utilized to determine the best-fit model. Among these models, a three break power-law provided the best fit to the XRT

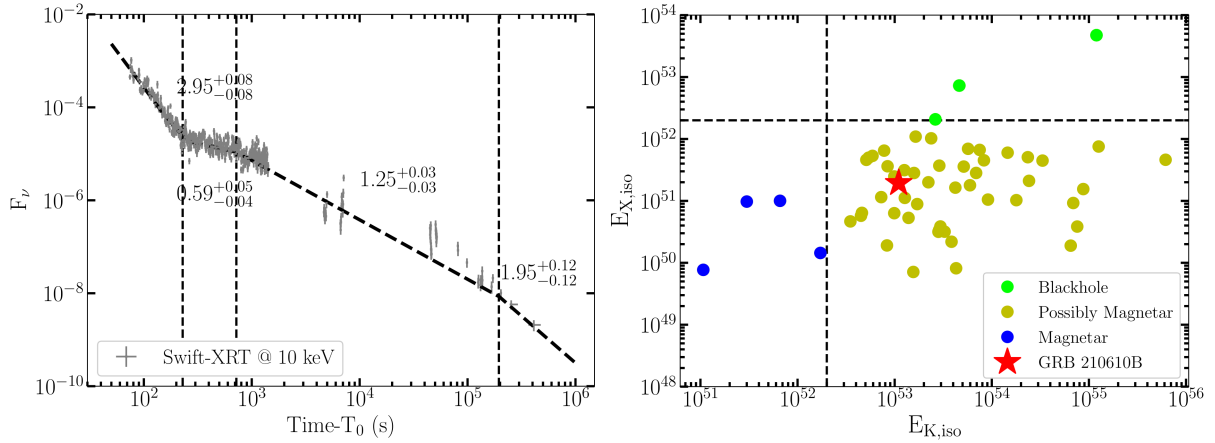


Figure 2: (Left) *Swift*-XRT flux density light curve of GRB 210610B. A three-break power-law fitted to the light curve is shown with a dashed line. The decay indices are mentioned on their corresponding slopes. The three vertical dashed lines represent the time corresponding to the breaks in the light curve. (Right) A comparison of observed energy released during the plateau phase ($E_{X,iso}$) vs. total kinetic energy release ($E_{K,iso}$) of the GRBs detected with *Swift*-XRT, taken from Ror et al. (in preparation).

light curve, and the obtained parameters are $\alpha_{x1} = 2.95^{+0.08}_{-0.08}$, $\alpha_{x2} = 0.59^{+0.05}_{-0.04}$, $\alpha_{x3} = 1.25^{+0.03}_{-0.03}$, $\alpha_{x4} = 1.95^{+0.12}_{-0.12}$, $t_{xb1} \simeq 220$ s, $t_{xb2} \simeq 700$ s, and $t_{xb3} \simeq 1.95 \times 10^5$ s. The *Swift*-XRT observation and the best-fit model curve are shown in Fig. 2.

3. Discussion and Conclusion

3.1. Origin of prompt emission

Time-resolved spectral analysis results of GRB 210610B show that α_{pt} is beyond the synchrotron line of death. However, Burgess et al. (2020) suggested that spectra can be well modelled with a synchrotron model even if the low-energy spectral index exceeds the synchrotron line-of-death. Indeed, we found that some of the bins are well fit by the synchrotron model. However, some of the bins favour the presence of superimposed thermal components as well. We studied the evolution of spectral parameters and found a rare feature where E_{pt} and B both showed flux-tracking behaviour throughout the prompt emission. The observed feature can be explained in terms of fireball cooling and expansion (Gupta et al., 2021; Ror et al., 2023). In the light of above, we suggest that GRB 210610B has a hybrid jet (Poynting flux outflow moving along with a hot fireball) composition which results in the synchrotron emission superimposed over a thermal component (Pe’er, 2015).

3.2. Progenitor

GRB emission is highly collimated. Therefore, the jet opening angle (θ_j) is a key parameter to get insights into its physics and energetics. One common method to estimate θ_j is from the fitting of the observed afterglow light curve. The time corresponding to the sudden fall in the X-ray light curve from the normal decay phase is considered as the jet break time T_j . The jet opening angle can be calculated (in radian) using the relation

$$\theta_j \simeq 0.057 \times T_{j,\text{days}}^{3/8} \times ((1+z)/2)^{3/8} \times (E_{\gamma,\text{iso},53})^{1/8} \times (\epsilon_{0.2})^{1/8} \times (n_{0.1})^{1/8},$$

where $T_{j,\text{days}}$ is the jet break time expressed in days, $E_{\gamma,\text{iso},53} = E_{\gamma,\text{iso}}/10^{53}$, $\epsilon_{0.2} = \epsilon/0.2$ and $n_{0.1} = n/0.1$ (n in cm^{-3}). GRBs with a T_{90} duration longer than 2 s are typically associated with the collapse of massive stars known as ‘‘collapsars.’’ However, recent discoveries by Ahumada et al. (2021) and Troja et al. (2022) have challenged this conventional understanding of the relationship between the T_{90} duration and GRB progenitors. To confirm whether the origin of GRB 210610B was indeed a collapsing massive star, we employed the relation presented in Bromberg et al. (2011):

$$T_{\text{Bore}}(\text{s}) \simeq 15 \times \epsilon_{\gamma}^{1/3} \times (L_{\gamma,\text{iso},50})^{1/3} \times (\theta_{10^\circ})^{2/3} \times (R_{11})^{2/3} \times (M_{15M_\odot})^{1/3},$$

where T_{Bore} represents the time required for the ultra-relativistic jet to penetrate the pre-existing cocoon surrounding the progenitor star, ϵ_{γ} is the radiative efficiency of the prompt emission, $L_{\gamma,\text{iso},50} = L_{\gamma,\text{iso}}/10^{50}$, $\theta_{10^\circ} = \theta_j/10^\circ$, while $R_{11} = R/10^{11}$ and $M_{15M_\odot} = M/(15M_\odot)$, R and M being the radius and the mass of the exploding star, resp. A method for the calculation of T_{Bore} given in Ror et al. (2023). $T_{90}/T_{\text{Bore}} > 1$ suggests that the burst originated from a collapsing massive star. For GRB 210610B, $T_{90}/T_{\text{Bore}} \simeq 107$, providing strong evidence that the most probable progenitor of this burst was indeed a collapsing massive star.

3.3. Central engine

One of the possible mechanisms for producing ultrarelativistic jets is through the formation of a highly rotating neutron star associated with extremely strong magnetic field lines (a millisecond magnetar). Pieces of evidence for a magnetar central engine are the existence of a plateau in the afterglow light curve and a highly polarized gamma-ray emission, associated with some GRBs, which could be produced by the strong magnetic field of a magnetar (Zhang and Mészáros, 2001). To constrain the central engine responsible for GRB 210610B, we employed two methods.

3.3.1. First method

It involved the calculation of beaming corrected energy release during the prompt emission. To obtain this, we multiplied the isotropic energy $E_{\gamma,\text{iso}}$ of the burst by a beaming correction factor defined as $f_b = 1 - \cos(\theta_j) \simeq 1/2(\theta_j)^2$, i.e., $E_{\gamma,\theta_j} = f_b \times E_{\gamma,\text{iso}}$. If the beaming corrected energy is greater than the maximum energy budget of a magnetar (i.e., $E_{\gamma,\theta_j} > 2 \times 10^{52}$ erg), then it ruled out the possibility of a magnetar central engine (Sharma et al., 2021). However, for GRB 210610B, the beaming-corrected energy is $E_{\gamma,\theta_j} = 1.06 \times 10^{51}$ erg $< 2 \times 10^{52}$ erg, suggesting that a magnetar could be the possible progenitor for this burst.

3.3.2. *Second method*

Under the assumption of synchrotron emission, the afterglow light curve of GRBs is expected to decay smoothly with a decay index of ~ 1 . However, some GRBs show a plateau in the afterglow light curve, which indicates that energy must be continuously supplied to the fireball to sustain the constant emission. A magnetar central is capable of providing such an energy injection (Zhang and Mészáros, 2001). However, several other possible scenarios can explain the observed plateau phase. In some cases, a moderately relativistic classical fireball is enough to explain the plateaus in the X-ray light curve caused by an external shock in a low-density wind-like surrounding medium (Dereli-Bégué et al., 2022). Further, GRB emission beamed narrowly in the forward direction with an opening angle $\theta_j \simeq 1/\Gamma$, with time, the Lorentz (Γ) factor of the jet decreases, and emission starts contributing from the off-axis region. This high-latitude emission can result in a plateau in the XRT light curve (Beniamini et al., 2020; Oganessian et al., 2020). Furthermore, Kumar et al. (2008) suggest that the continued accretion due to small viscous parameters and fall-back of residual gas on the central engine can also cause the plateaus in the X-ray light curve.

We have calculated the X-ray isotropic energy release, $E_{X,\text{iso}}$, using the relation given in Li et al. (2018). A comparison between the energy released during the plateau phase ($E_{X,\text{iso}}$) and the total kinetic energy release ($E_{K,\text{iso}}$) of the GRBs detected by the *Swift*-XRT instrument is shown in Fig. 2 taken from Ror et al. (in preparation). For GRB 210610B, the obtained value of $E_{X,\text{iso}} = 1.87 \times 10^{51}$ erg, which is less than the total energy budget of a magnetar, once again favoring the magnetar central engine as the likely progenitor for this burst. The combined energy emitted during the prompt+afterglow emission phase is 2.94×10^{51} erg $< 2 \times 10^{52}$ erg favoring the magnetar central engine.

4. Future Prospect

This article examines the discernible characteristics of GRB 210610B through the utilization of archival data from space-based observations. Our forthcoming objective involves extending the analysis into the realm of multi-wavelength observations. Additionally, we intend to compare the afterglow emission of GRB 210610B with a collection of light curves from similar bursts to gain deeper insights into the burst's underlying progenitors, central engine, and surrounding environment.

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Author contributions

RG has contributed to the prompt and afterglow analyses of the burst. SBP and AA thoroughly reviewed the manuscript and provided insightful comments that greatly improved the quality of the draft.

Conflicts of interest

The authors affirm that they do not have any conflicts of interest regarding the manuscript.

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